

**Explicit/Implicit, Conservative,
Galerkin Domain Decomposition
Procedures for Parabolic Problems**

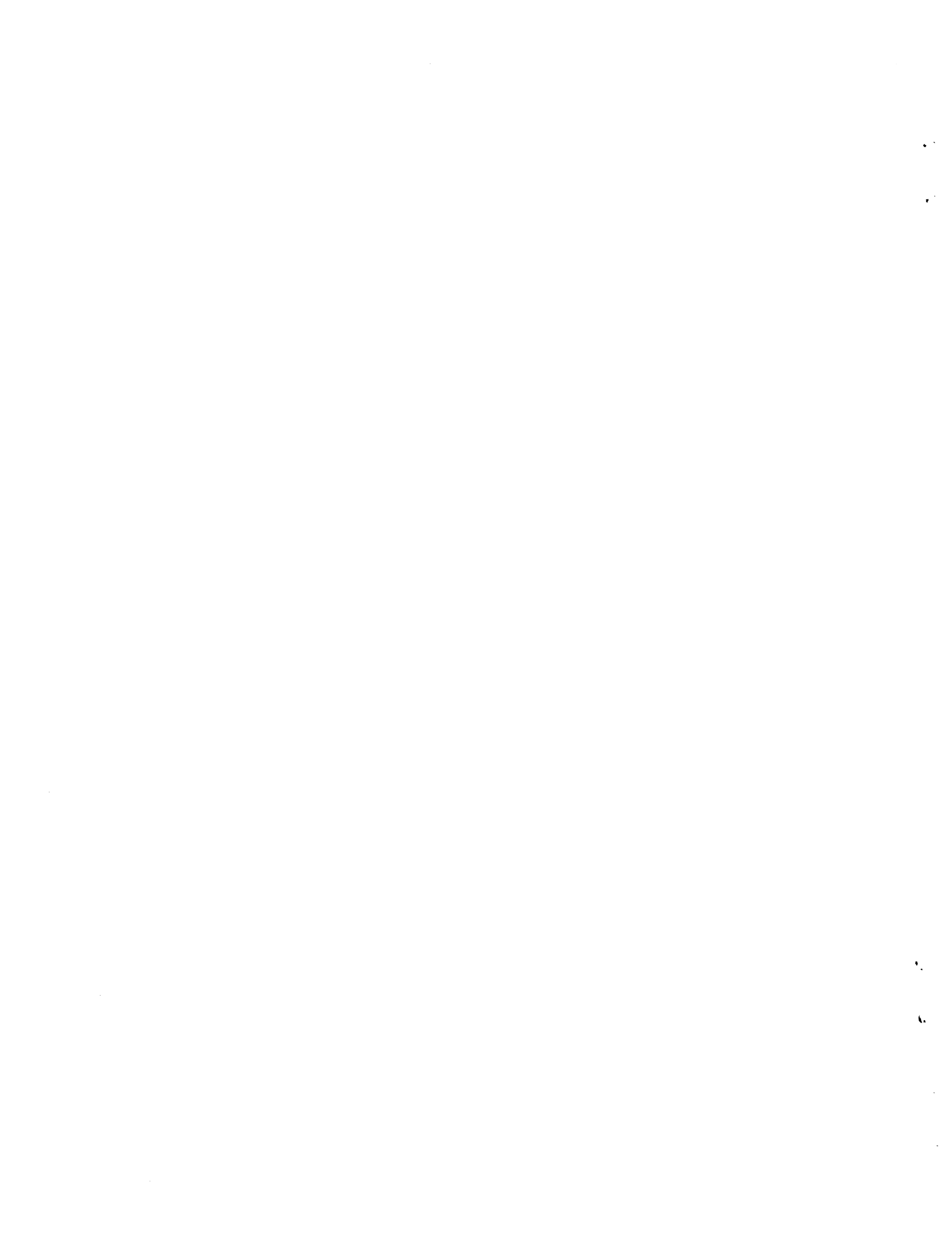
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Explicit/Implicit, Conservative, Galerkin Domain Decomposition Procedures for Parabolic Problems

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Abstract: Several domain decomposition methods for approximating solutions of parabolic problems are given. These methods rely on implicit Galerkin procedures in the subdomains and explicit flux calculation on the inter-domain boundaries. The procedures are conservative both in the subdomains and across inter-domain boundaries. *A priori* error bounds and experimental results are presented.

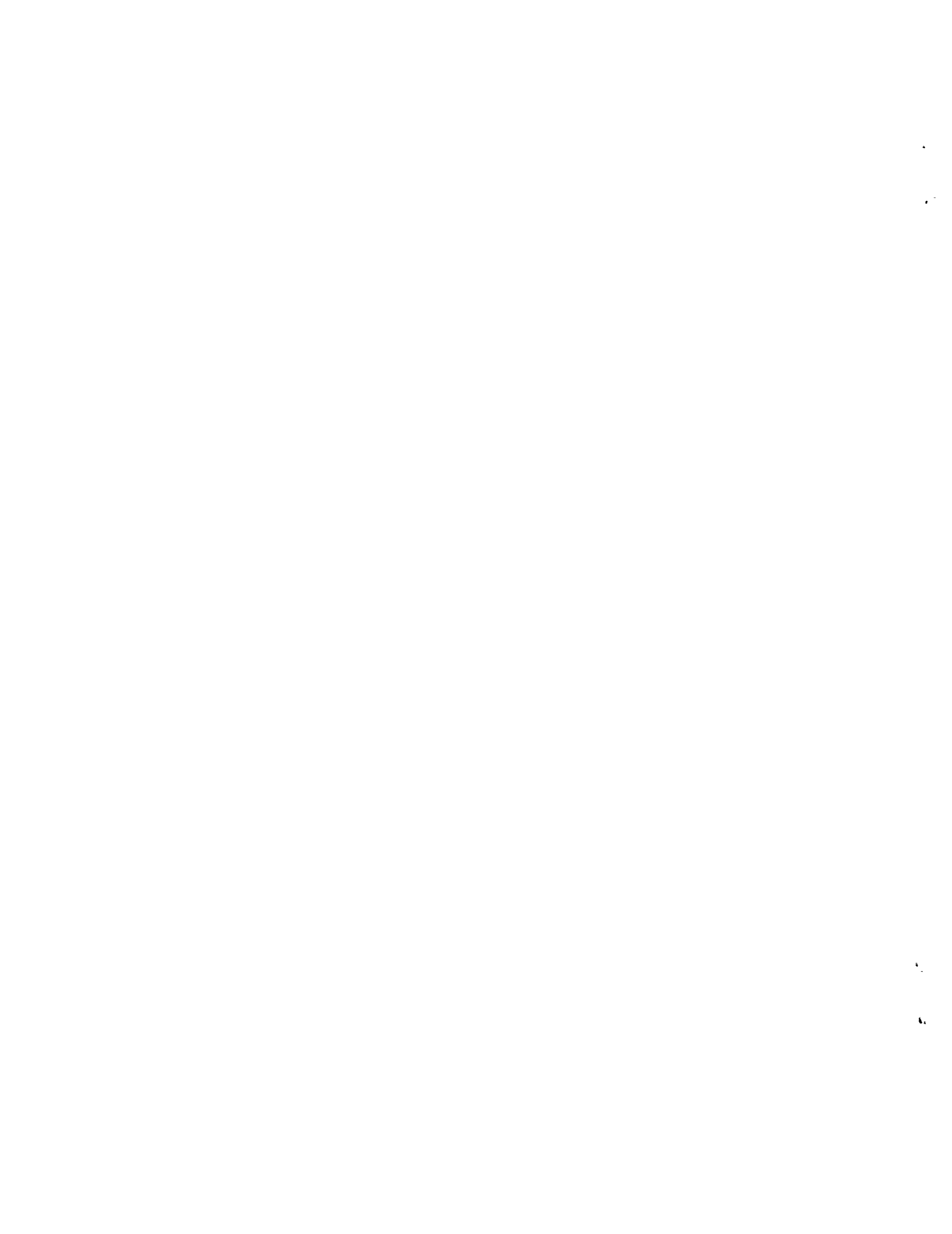
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1 Introduction

When solving parabolic partial differential equations using Galerkin finite element procedures, it is frequently advantageous to use implicit time-stepping, due to the severe time-step constraint needed to insure stability when using explicit time-stepping. Thus, a large, global system of equations must be solved at each time step. Domain decomposition procedures can be used to break these large computations into several smaller ones, but the smaller

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“subdomain problems” must be coupled in some way. The procedures studied here use simple, explicit calculations on the boundaries between subdomains to predict the flux and this is the only coupling between subproblems. Thus, these procedures are non-iterative, and involve dividing the domain into nonoverlapping subdomains. The explicit nature of the flux prediction induces a time step limitation that is necessary to preserve stability, but this constraint is less severe than that which comes with a fully explicit method.

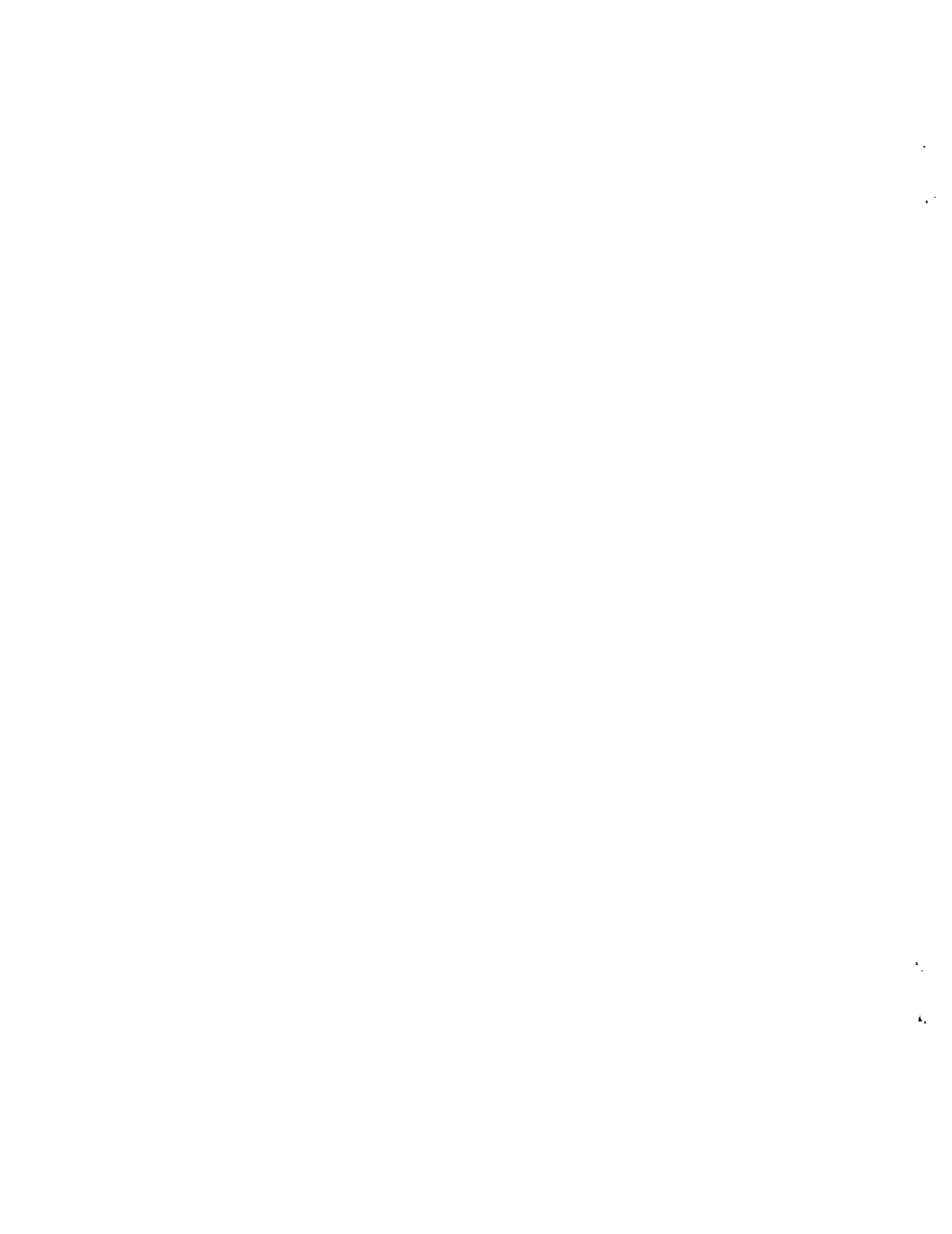
Galerkin procedures are useful in situations that require non-rectangular geometry and the procedures given here allow considerable geometric flexibility. The function spaces used on the subdomains need not match up in such a way that they are restrictions of the same global subspace of H^1 , and the operator used to approximate the flux does not have to match the grids used in the subdomains.

The *a priori* error bounds given here rely on previous work on elliptic Galerkin approximations. In fact, our theorems are stated in terms of the errors of certain elliptic approximations rather than powers of some asymptotic parameter.

We exhibit second and fourth order correct boundary flux approximations. The analysis is done for arbitrary order correct boundary flux approximations. The first results use a first order time discretization, but the results in Section 4 give a second order in time approximation.

Experimental results are also given. They show that the stability constraint is necessary, although not necessarily sharp. They also exhibit (on some example problems) asymptotic rates of convergence that are slightly better than are proved here. These experimental rates of convergence agree with what can be proved in some special cases.

In a previous paper on finite differences [3], we coupled subdomains by explicit predictions of boundary values and this work was carried over to a finite element context in



[2]. The new procedures given here offer greater geometric generality. They also have the advantage that they are conservative across subdomains boundaries in the sense that the interior boundaries do not serve as sources or sinks. This conservation property is missing in our earlier schemes based on explicit interface value prediction.

2 Preliminaries

Let Ω denote a spatial domain in \mathbf{R}^d . Denote by $H^m(\Omega)$ and $W_\infty^m(\Omega)$ the standard Sobolev spaces on Ω , with norms $\|\cdot\|_m$ and $\|\cdot\|_{\infty,m}$, respectively. Let $L^p(\Omega)$, $p = 2, \infty$, denote the standard Banach spaces, with $\|\cdot\|$ denoting the L^2 norm, $\|\cdot\|_\infty$ the L^∞ norm.

Let $[\alpha, \beta] \subset [0, T]$ denote a time interval, $X = X(\Omega)$ a normed space. To incorporate time dependence, we use the notation $\|\cdot\|_{L^p(\alpha, \beta; X)}$ to denote the norm of X -valued functions f with the map $t \mapsto \|f(\cdot, t)\|_X$ belonging to $L^p(\alpha, \beta)$.

Assume Ω has a piecewise uniformly smooth, Lipschitz boundary, $\partial\Omega$. Assume that u^0 , a , and b are smooth, real-valued functions on $\bar{\Omega}$, with a being positive and b nonnegative. For some $T > 0$, the function $u(x, t)$ satisfies

$$\frac{\partial u}{\partial t} - \nabla \cdot (a \nabla u) + bu = 0, \quad \text{on } \Omega \times (0, T], \quad (2.1)$$

$$\frac{\partial u}{\partial n_\Omega} = 0, \quad \text{on } \partial\Omega \times (0, T], \quad (2.2)$$

$$u(x, 0) = u^0(x), \quad \text{on } \Omega, \quad (2.3)$$

where n_Ω is the outward normal to $\partial\Omega$.

We use approximations of derivatives of delta functions at several points in this work, and these approximations can be viewed as coming from one-dimensional approximations of the delta function. For future reference we define two special functions ϕ_2 and ϕ_4 as

follows:

$$\phi_2(x) = \begin{cases} 1 - x, & 0 \leq x \leq 1, \\ x + 1, & -1 \leq x \leq 0, \\ 0, & \text{otherwise.} \end{cases}$$

$$\phi_4(x) = \begin{cases} (x - 2)/12, & 1 \leq x \leq 2, \\ -5x/4 + 7/6, & 0 \leq x \leq 1, \\ 5x/4 + 7/6, & -1 \leq x \leq 0, \\ -(x + 2)/12, & -2 \leq x \leq -1, \\ 0, & \text{otherwise.} \end{cases}$$

Note that if $p(x)$ is a polynomial of degree at most one, then

$$\int p(x)\phi_2(x)dx = p(0), \quad (2.4)$$

and if p is a polynomial of degree at most three,

$$\int p(x)\phi_4(x)dx = p(0).$$

3 Domain Decomposition Procedures

In this section we consider the case of two subdomains and we use a Galerkin procedure on each subdomain. As motivation for the abstract formulation presented later we start with an example. For this example use $a \equiv b \equiv 1$ in (2.1).

Let $\Omega = (0, 1) \times (0, 1)$. Take $\Omega_1 = (0, \frac{1}{2}) \times (0, 1)$, $\Omega_2 = (\frac{1}{2}, 1) \times (0, 1)$, and $\Gamma = \{\frac{1}{2}\} \times (0, 1)$.

Define, for some $H \in (0, \frac{1}{2})$,

$$\phi(x) = \phi_2\left(\left(x - \frac{1}{2}\right)/H\right)/H. \quad (3.1)$$

For $j = 1, 2$, let \mathcal{M}_j be a finite dimensional subspace of $H^1(\Omega_j)$, and let \mathcal{M} be the subspace of $L^2(\Omega)$ such that if $v \in \mathcal{M}$, $v|_{\Omega_j} \in \mathcal{M}_j$. Note that functions v in \mathcal{M} have a well-defined jump $[v]$ on Γ ; $[v](\frac{1}{2}, y) = v(\frac{1}{2}, y + 0) - v(\frac{1}{2}, y - 0)$.

Take the bilinear form D to be given by

$$D(\psi, \rho) = \sum_{i=1}^2 \int_{\Omega_i} (\nabla \psi \cdot \nabla \rho + \psi \rho) dx dy. \quad (3.2)$$

Define an approximate derivative as follows:

$$B(\psi) \left(\frac{1}{2}, y \right) = - \int_0^1 \phi'(x) \psi(x, y) dx, \quad (3.3)$$

where ϕ is given by (3.1).

Let $(\cdot, \cdot)_{\Sigma}$ be the $L^2(\Sigma)$ inner product:

$$(\psi, \rho)_{\Sigma} = \int_{\Sigma} \psi \rho. \quad (3.4)$$

In the case $\Sigma = \Omega$ we omit the subscript:

$$(\psi, \rho) = (\psi, \rho)_{\Omega}. \quad (3.5)$$

Let $0 = t^0 < t^1 < \dots < t^M = T$ be a given sequence and suppose that $U^0 \in \mathcal{M}$ is given.

Define U^1, \dots, U^M by

$$(\partial_t U^n, v) + D(U^n, v) + (B(U^{n-1}), [v])_{\Gamma} = 0, \quad v \in \mathcal{M}, \quad (3.6)$$

where $\Delta t^n = t^n - t^{n-1}$, and

$$\partial_t U^n = (U^n - U^{n-1}) / \Delta t^n. \quad (3.7)$$

This scheme has the property that U^n can be computed on Ω_1 and Ω_2 completely independently once $B(U^{n-1})$ has been computed on Γ . The flux at each point on Γ is computed explicitly from U^{n-1} , and the two parts of U^n are then computed using an implicit Galerkin

backward difference equation on each Ω_j . Note also that if the function $v \equiv 1$ belongs to \mathcal{M}_j , then (3.6) is conservative in the following sense. If we use $b \equiv 0$ instead of 1, then the average value of U^n is the same for all n , just as the average value of $u(\cdot, t)$ is independent of t , by (2.1) and (2.2).

For functions ψ with restrictions in $H^1(\Omega_1)$ and $H^1(\Omega_2)$ define

$$|||\psi|||^2 = D(\psi, \psi) + H^{-1} ||[\psi]||_{L^2(\Gamma)}^2. \quad (3.8)$$

Next, note that for such ψ 's,

$$\begin{aligned} (B(\psi), [\psi])_\Gamma &= - \int_0^1 \int_0^1 \phi'(x) \psi(x, y) [\psi] \left(\frac{1}{2}, y \right) dx dy \\ &= \frac{1}{H} \int_0^1 [\psi]^2 \left(\frac{1}{2}, y \right) dy \\ &\quad + \int_0^1 \int_0^1 \phi(x) \psi_x(x, y) dx [\psi] \left(\frac{1}{2}, y \right) dy, \end{aligned} \quad (3.9)$$

and

$$\begin{aligned} \int_0^1 \int_0^1 \phi(x) \psi_x(x, y) dx [\psi] \left(\frac{1}{2}, y \right) dy \\ \leq \|\psi_x\|_{L^2(\Omega_1 \cup \Omega_2)} \|\phi\|_{L^2(0,1)} ||[\psi]||_{L^2(\Gamma)} \\ \leq \|\psi_x\|_{L^2(\Omega_1 \cup \Omega_2)} \left(\left(\frac{2}{3H} \right)^{\frac{1}{2}} ||[\psi]||_{L^2(\Gamma)} \right). \end{aligned} \quad (3.10)$$

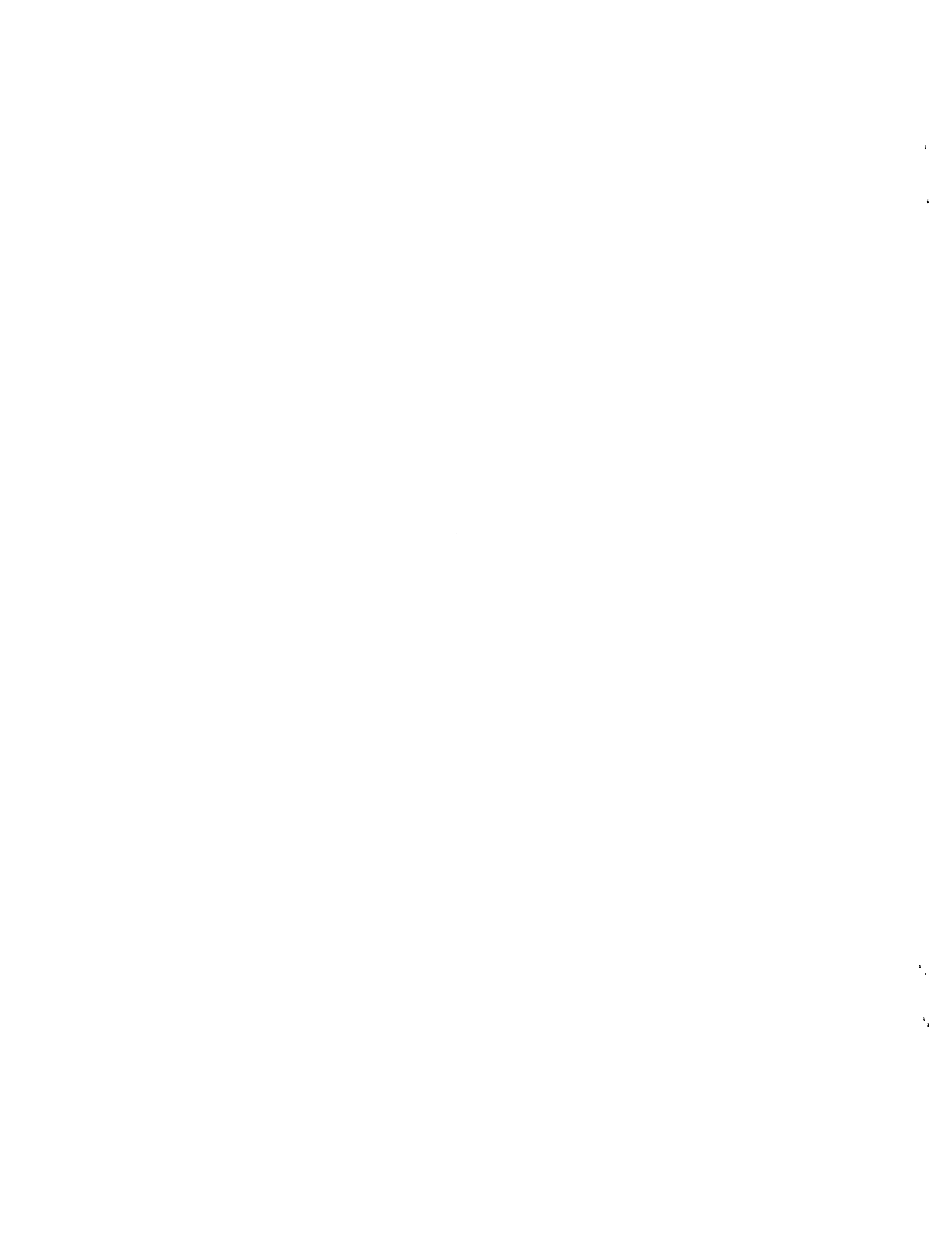
From this result we get that

$$D(\psi, \psi) + (B(\psi), [\psi])_\Gamma \geq \left(1 - \frac{1}{\sqrt{6}} \right) |||\psi|||^2; \quad (3.11)$$

in deriving this we used

$$\alpha\beta \leq \epsilon\alpha^2 + \frac{1}{4\epsilon}\beta^2 \quad (3.12)$$

on the product in (3.10) with $\epsilon^2 = 1/6$.



We restate (3.11) as

$$\begin{aligned} |||\psi|||^2 &\leq \left(1 - \frac{1}{\sqrt{6}}\right)^{-1} (D(\psi, \psi) + (B(\psi), [\psi])_\Gamma) \\ &\leq 1.7 (D(\psi, \psi) + (B(\psi), [\psi])_\Gamma). \end{aligned} \quad (3.13)$$

The following two bounds are straightforward

$$\|B(\psi)\|_{L^2(\Gamma)}^2 \leq 2H^{-3}\|\psi\|^2, \quad (3.14)$$

$$\|B(\psi)\|_{L^2(\Gamma)} \leq 2H^{-1}\|\psi\|_\infty. \quad (3.15)$$

Also, if ψ is smooth in Ω , it follows from integration by parts, (2.4), and Taylor's Theorem that

$$B(\psi) \left(\frac{1}{2}, y\right) = \psi_x \left(\frac{1}{2}, y\right) + \int_0^1 \int_{\frac{1}{2}}^x (x-s) \psi_{xxx} \left(\frac{1}{2}, s\right) ds \phi(x) dx. \quad (3.16)$$

Hence it follows

$$\left\| \psi_x \left(\frac{1}{2}, \cdot\right) - B(\psi) \right\|_{L^2(\Gamma)} \leq \frac{H^2}{12} \|\psi_{xxx}\|_\infty. \quad (3.17)$$

To state an error bound for the procedure (3.6), we introduce an elliptic projection [6] $W \in \mathcal{M}$ of the solution u as follows:

$$D(u(\cdot, t) - W(\cdot, t), v) = 0, \quad v \in \mathcal{M}. \quad (3.18)$$

The function W at each t is just the $H^1(\Omega_j)$ projection of u into \mathcal{M}_j . Let the error in the projection be denoted by

$$\eta = u - W. \quad (3.19)$$

Theorem 1 *Suppose that the solution u is sufficiently smooth and that $U^0 \in \mathcal{M}$ is taken to be W^0 . Let $\Delta t = \max_n \Delta t^n$. Then there exists a constant C , independent of the spaces*

\mathcal{M}_j , such that

$$\max_n \|u^n - U^n\| \leq C \left(\Delta t + H^{2.5} + \int_0^T \|\eta_t(\cdot, t)\| dt + H^{-\frac{1}{2}} \|\eta\|_{L^\infty(\Omega \times (0, T))} \right), \quad (3.20)$$

provided that

$$\Delta t \leq \frac{H^2}{4}. \quad (3.21)$$

The procedure is first order correct in Δt as expected since this would be the case even if no domain decomposition were used. The second order correct approximation B gives rise to the $H^{2.5}$ since we only make the H^2 error on a “small” set. The loss of $H^{-\frac{1}{2}}$ in the last term in the estimate can be avoided in certain special cases [1], but at this level of generality we do not know how to improve this term. The theorem will follow as an easy consequence of Theorem 2 below.

We now treat a more general case. The domain Ω in \mathbf{R}^d is divided into two nonoverlapping subdomains Ω_1 and Ω_2 . The interface between these domains

$$\Gamma = \Omega \cap \bar{\Omega}_1 \cap \bar{\Omega}_2 \quad (3.22)$$

is assumed to be a uniformly smooth, $(d - 1)$ -dimensional manifold.

For functions ψ, ρ with restrictions in $H^1(\Omega_j)$ let

$$D(\psi, \rho) = \int_{\Omega_1 \cup \Omega_2} (a \nabla \psi \cdot \nabla \rho + b \psi \rho) dx, \quad (3.23)$$

and let the jump in such a ψ on Γ be denoted by $[\psi]$. For definiteness let $[\psi]$ be the trace from Ω_2 minus the trace from Ω_1 .

We need a parametrized approximation to the normal derivative on Γ . Assume for some $H > 0$ that B is a linear map of $L^2(\Omega)$ into $L^2(\Gamma)$ and that it satisfies the following four conditions, which are generalizations of (3.13), (3.14), (3.15), and (3.17):

(i) There is a constant C_0 such that

$$|||\psi|||^2 \leq C_0 (D(\psi, \psi) + (aB(\psi), [\psi])_\Gamma), \quad (3.24)$$

where

$$|||\psi|||^2 = D(\psi, \psi) + \frac{1}{H} (a[\psi], [\psi])_\Gamma. \quad (3.25)$$

(ii) There is a constant C_1 such that

$$(aB(\psi), B(\psi))_\Gamma \leq C_1 H^{-3} |||\psi|||^2. \quad (3.26)$$

(iii) There is a constant C_2 such that

$$\|B(\psi)\|_{L^2(\Gamma)} \leq C_2 H^{-1} \|\psi\|_{L^\infty(\Omega)}. \quad (3.27)$$

(iv) There is a constant $k \geq 0$ and a constant C_3 which depends on the solution u such that

$$\left\| \frac{\partial u(\cdot, t)}{\partial \gamma} - B(u)(\cdot, t) \right\|_{L^2(\Gamma)} \leq C_3 H^k, \quad (3.28)$$

for $0 \leq t \leq T$, where $\partial u / \partial \gamma$ is the normal derivative of u on Γ , in the direction from Ω_1 to Ω_2 .

Suppose that \mathcal{M}_j is a finite dimensional subspace of $H^1(\Omega_j)$ and let \mathcal{M} be the set of $L^2(\Omega)$ functions whose restrictions to Ω_j belong to \mathcal{M}_j . For $0 = t^0 < t^1 < \dots < t^M = T$, and $U^0 \in \mathcal{M}$ given, define U^1, \dots, U^M by

$$(\partial_t U^n, v) + D(U^n, v) + (aB(U^{n-1}), [v])_\Gamma = 0, \quad v \in \mathcal{M}, \quad (3.29)$$

where $\Delta t^n = t^n - t^{n-1}$,

$$\partial_t U^n = (U^n - U^{n-1}) / \Delta t^n. \quad (3.30)$$

Just as in the example we define $W(\cdot, t) \in \mathcal{M}$ by

$$D((u - W)(\cdot, t), v) = 0, \quad v \in \mathcal{M}, \quad (3.31)$$

at each $t \in [0, T]$. Let the error in the elliptic projection be denoted by

$$\eta = u - W. \quad (3.32)$$

Theorem 2 *Suppose that the solution u to (2.1) is sufficiently smooth and that $U^0 = W(\cdot, 0)$. Let $\Delta t = \max_n \Delta t^n$. Then there exists a constant C , independent of the spaces \mathcal{M}_j such that*

$$\max_n \|(u - U)(\cdot, t^n)\| \leq C \left\{ \Delta t + H^{k+\frac{1}{2}} + \int_0^T \|\eta_t(\cdot, t)\| dt + H^{-\frac{1}{2}} \|\eta\|_{L^\infty(\Omega \times (0, T))} \right\},$$

provided

$$\Delta t \leq \frac{H^2}{C_0 C_1}. \quad (3.33)$$

Note that there are no assumptions that require \mathcal{M}_1 and \mathcal{M}_2 to be compatible in some way. Also, as we saw in the example, the H parameter for the operator B is not necessarily related to any aspect of the spaces \mathcal{M}_j ; in particular, it is not required that ϕ restricted to Ω_j have any relation to \mathcal{M}_j .

Proof of Theorem 2. Note first that u satisfies

$$(\partial_t u^n, v) + D(u^n, v) + \left(a \frac{\partial u^n}{\partial \gamma}, [v] \right)_\Gamma = (\rho^n, v), \quad v \in \mathcal{M}, \quad (3.34)$$

where $u^n(x) = u(x, t^n)$. The time truncation term ρ^n satisfies

$$\begin{aligned} \sum_{n=1}^M \|\rho^n\| \Delta t^n &\leq \Delta t \int_0^T \|u_{tt}(\cdot, t)\| dt \\ &\leq C \Delta t. \end{aligned} \quad (3.35)$$

From (3.34) and the definition of W we see that

$$\begin{aligned} & (\partial_t W^n, v) + D(W^n, v) + (aB(W^{n-1}), [v])_\Gamma \\ &= (\rho^n - \partial_t \eta^n, v) + \left(a \left(B(W^{n-1}) - \frac{\partial u^n}{\partial \gamma} \right), [v] \right)_\Gamma, \quad v \in \mathcal{M}. \end{aligned} \quad (3.36)$$

Let

$$\nu = U - W. \quad (3.37)$$

This gives, from (3.29) and (3.36),

$$\begin{aligned} & (\partial_t \nu^n, v) + D(\nu^n, v) + (aB(\nu^{n-1}), [v])_\Gamma \\ &= (\partial_t \eta^n - \rho^n, v) + \left(a \left(\frac{\partial u^n}{\partial \gamma} - B(W^{n-1}) \right), [v] \right)_\Gamma, \quad v \in \mathcal{M}. \end{aligned} \quad (3.38)$$

Let

$$\vartheta = a \left(\frac{\partial u^n}{\partial \gamma} - B(W^{n-1}) \right). \quad (3.39)$$

Then

$$\vartheta = a \left(\frac{\partial u^n}{\partial \gamma} - \frac{\partial u^{n-1}}{\partial \gamma} + \frac{\partial u^{n-1}}{\partial \gamma} - B(u^{n-1}) + B(\eta^{n-1}) \right).$$

From this it follows that

$$\|\vartheta\|_{L^2(\Gamma)} \leq C(\Delta t + H^k + H^{-1}\|\eta^{n-1}\|_\infty). \quad (3.40)$$

Use $v = \nu^n$ in (3.38) to see that

$$\begin{aligned} & (\partial_t \nu^n, \nu^n) + D(\nu^n, \nu^n) + (aB(\nu^n), [\nu^n])_\Gamma \\ & \leq (aB(\nu^n - \nu^{n-1}), [\nu^n])_\Gamma + \|\partial_t \eta^n + \rho^n\| \|\nu^n\| \\ & \quad + C \left(\Delta t + H^k + H^{-1}\|\eta^{n-1}\|_\infty \right) H^{\frac{1}{2}} \|\nu^n\|. \end{aligned} \quad (3.41)$$

Since $(\alpha - \beta)\alpha = \frac{1}{2}(\alpha^2 - \beta^2) + \frac{1}{2}(\alpha - \beta)^2$, we see that

$$(\partial_t \nu^n, \nu^n) = \frac{1}{2} \partial_t (\|\nu^n\|^2) + \frac{\Delta t^n}{2} \|\partial_t \nu^n\|^2. \quad (3.42)$$

Use this and (3.24) in (3.41) to see that

$$\begin{aligned} & \partial_t (\|\nu^n\|^2) + \Delta t^n \|\partial_t \nu^n\|^2 + \frac{2}{C_0} \|\nu^n\|^2 \\ & \leq 2C_1^{\frac{1}{2}} H^{-1} \Delta t^n \|\partial_t \nu^n\| \|\nu^n\| + 2\|\partial_t \eta^n + \rho^n\| \|\nu^n\| \\ & \quad + C(\Delta t + H^k + H^{-1} \|\eta^{n-1}\|_\infty) H^{\frac{1}{2}} \|\nu^n\| \\ & \leq \frac{2}{C_0} \|\nu^n\|^2 + (C_0 C_1 H^{-2} \Delta t^n) \Delta t^n \|\partial_t \nu^n\|^2 \\ & \quad + 2\|\partial_t \eta^n + \rho^n\| \|\nu^n\| + C(\Delta t + H^k + H^{-1} \|\eta^{n-1}\|_\infty)^2 H. \end{aligned} \quad (3.43)$$

Provided

$$\Delta t^n \leq \frac{H^2}{C_0 C_1}, \quad (3.44)$$

we see that

$$\partial_t (\|\nu^n\|^2) \leq 2\|\partial_t \eta^n + \rho^n\| \|\nu^n\| + C(\Delta t^2 H + H^{2k+1} + H^{-1} \|\eta^{n-1}\|_\infty^2). \quad (3.45)$$

From this we see that, by multiplying by Δt^n and summing on n ,

$$\begin{aligned} (\max_n \|\nu^n\|)^2 & \leq \|\nu^0\|^2 + 2 \sum_{n=1}^M \|\partial_t \eta^n + \rho^n\| \|\nu^n\| \Delta t^n \\ & \quad + C(\Delta t^2 H + H^{2k+1} + H^{-1} \|\eta\|_{L^\infty(\Omega \times (0, T))}^2). \end{aligned} \quad (3.46)$$

Use $\nu^0 = 0$ and

$$\begin{aligned} 2 \sum_{n=1}^M \|\partial_t \eta^n + \rho^n\| \|\nu^n\| \Delta t^n & \leq 2 \max_n \|\nu^n\| \sum_{n=1}^M \|\partial_t \eta^n + \rho^n\| \Delta t^n \\ & \leq \frac{1}{2} (\max_n \|\nu^n\|)^2 + 4 \left(\sum_{n=1}^M \|\partial_t \eta^n + \rho^n\| \Delta t^n \right)^2 \end{aligned} \quad (3.47)$$

to see that

$$\begin{aligned} \max_n \|\nu^n\| \leq C \left(\Delta t + H^{k+\frac{1}{2}} + H^{-\frac{1}{2}} \|\eta\|_{L^\infty(\Omega \times (0,T))} \right. \\ \left. + \int_0^T \|\eta_t(\cdot, t)\| dt \right). \end{aligned} \quad (3.48)$$

Here we have also used (3.35). The conclusion follows from the triangle inequality. //

The bound (3.44) (and hence (3.33)) can be relaxed by almost a factor of two, at the expense of a larger constant in the last term in (3.43). The L^∞ bound on η is needed near Γ in the examples of B with which we have worked; away from Γ an L^2 bound can be used instead.

Note that the projection W has no dependence on H , and this implies that η has no H dependence. From this observation it follows that we can allow H to vary from step to step, provided only that

$$\Delta t^n \leq \frac{H_n^2}{C_0 C_1}. \quad (3.49)$$

Return to consideration of the example problem used to introduce this section. Define B using ϕ_4 of Section 2 by taking $\phi(x) = \phi_4((x - \frac{1}{2})/H)/H$. Now we need $H \in (0, \frac{1}{4})$. Then calculation gives that (3.24), (3.26), and (3.27) hold with $C_0 = 1.64$, $C_1 = 3.14$, and $C_2 = 8/3$. Thus, Theorem 2 requires

$$\Delta t \leq \frac{H^2}{5.15}.$$

In this case (3.28) holds with $k = 4$.

4 A Second Order in Time Procedure

In this section we illustrate the use of a second order in time backward discretization using the Galerkin based procedure of Section 3. We restrict attention to the case of uniform

time steps $t^n = n\Delta t$. The context used here is that of Theorem 2; i.e., we allow variable coefficients, $\Omega \subset \mathbb{R}^d$, and Γ is a smooth $(d - 1)$ -dimensional manifold.

The results of this section are slightly related to the results in [5], where energy methods are used on second order backward difference Galerkin methods (and blending methods). In [5], variable time steps are analyzed, but not in a domain decomposition context.

Let

$$\delta_2 U^n = \partial_t U^n + \frac{1}{2}(\partial_t U^n - \partial_t U^{n-1}). \quad (4.1)$$

Assume that U^0 and U^1 in \mathcal{M} are known and define U^2, \dots, U^M by the following analogue of (3.29):

$$(\delta_2 U^n, v) + D(U^n, v) + \left(B(2U^{n-1} - U^{n-2}), [v] \right)_\Gamma = 0, \quad v \in \mathcal{M}. \quad (4.2)$$

This procedure is formally second order correct in time, and like (3.29), it allows the calculation of U^n to be done independently on Ω_1 and Ω_2 once the operator B has been evaluated on Γ using $2U^{n-1} - U^{n-2}$.

The analogue of Theorem 2 is the following:

Theorem 3 *Take W and η as in (3.31) and (3.32) and let B satisfy (3.24), (3.26), (3.27), and (3.28). Suppose that the solution u of (2.1) is sufficiently smooth and that $U^0 = W(\cdot, 0)$ and $U^1 = W(\cdot, \Delta t)$. Then there is a constant C , independent of the spaces \mathcal{M}_j , such that U^n given by (4.2) satisfies*

$$\max_n \|(u - U)(\cdot, t^n)\| \leq C \left\{ \Delta t^2 + H^{k+1/2} + \int_0^T \|\eta_t(\cdot, t)\| dt + H^{-1/2} \|\eta\|_{L^\infty(\Omega \times (0, T))} \right\},$$

provided

$$\Delta t \leq \frac{H^2}{2C_0 C_1}. \quad (4.3)$$

Of course it is unlikely that we could take U^0 and U^1 exactly as indicated in the theorem. The proof will show that U^0 and U^1 need only approximate these values in a certain sense, but we will not give the details of that result.

Proof of Theorem 3. In analogy with (3.34) we see that

$$(\delta_2 u^n, v) + D(u^n, v) + \left(a \frac{\partial u^n}{\partial \gamma}, [v] \right)_\Gamma = (\rho^n, v), \quad v \in \mathcal{M}, \quad (4.4)$$

where

$$\begin{aligned} \sum_{n=1}^M \|\rho^n\| \Delta t &\leq C(\Delta t)^2 \int_0^T \|u_{tt}(\cdot, t)\| dt \\ &\leq C(\Delta t)^2. \end{aligned} \quad (4.5)$$

Thus

$$\begin{aligned} &(\delta_2 W^n, v) + D(W^n, v) + \left(aB(2W^{n-1} - W^{n-2}), [v] \right)_\Gamma \\ &= (\rho^n - \delta_2 \eta^n, v) + \left(a \left(B(2W^{n-1} - W^{n-2}) - \frac{\partial u^n}{\partial \gamma} \right), [v] \right)_\Gamma, \quad v \in \mathcal{M}_j. \end{aligned} \quad (4.6)$$

Use

$$\begin{aligned} &\frac{\partial u^n}{\partial \gamma} - B(2W^{n-1} - W^{n-2}) \\ &= \frac{\partial(u^n - 2u^{n-1} + u^{n-2})}{\partial \gamma} + \left[\frac{\partial(2u^{n-1} - u^{n-2})}{\partial \gamma} - B(2u^{n-1} - u^{n-2}) \right] \\ &\quad + B(2\eta^{n-1} - \eta^{n-2}) \end{aligned} \quad (4.7)$$

to get

$$\begin{aligned} &(\delta_2 \nu^n, \nu^n) + D(\nu^n, \nu^n) + (aB(\nu^n), [\nu^n])_\Gamma \\ &= \left(aB(\nu^n - 2\nu^{n-1} + \nu^{n-2}), [\nu^n] \right)_\Gamma + \|\delta_2 \eta^n + \rho^n\| \|\nu^n\| \\ &\quad + C \left((\Delta t)^2 + H^k + H^{-1}(\|\eta^{n-1}\|_\infty + \|\eta^{n-2}\|_\infty) \right) H^{1/2} \|\nu^n\|. \end{aligned} \quad (4.8)$$

To replace (3.42) we use

$$(\delta_2 \nu^n, \nu^n) = \partial_t E(\nu^n, \nu^{n-1}) + \frac{1}{4\Delta t} (\nu^n - 2\nu^{n-1} + \nu^{n-2})^2 \quad (4.9)$$

where

$$\begin{aligned} E(\psi, \mu) &= \frac{1}{4} (\|\psi\|^2 + \|2\psi - \mu\|^2) \\ &\geq \frac{1}{4} \|\psi\|^2. \end{aligned} \quad (4.10)$$

The first term on the right side of (4.8) is bounded by

$$\begin{aligned} &\frac{C_0 H}{2} \left(aB(\nu^n - 2\nu^{n-1} + \nu^{n-2}), B(\nu^n - 2\nu^{n-1} + \nu^{n-2}) \right)_\Gamma + \frac{1}{2C_0} \|\nu^n\|^2 \\ &\leq \frac{C_0 H}{2} C_1 H^{-3} \|\nu^n - 2\nu^{n-1} + \nu^{n-2}\|^2 + \frac{1}{2C_0} \|\nu^n\|^2, \end{aligned}$$

where (3.26) was used to introduce C_1 . To complete the proof follow the proof of Theorem

2. //

The critical step in determining the Δt constraint is the relation

$$\frac{1}{4\Delta t} - \left(\frac{C_0 H}{2} \right) \left(\frac{C_1}{H^3} \right) \geq 0.$$

This gives

$$\Delta t \leq \frac{H^2}{2C_0 C_1}; \quad (4.11)$$

just as in Theorem 2 this could be relaxed by almost a factor of two.

5 Numerical experiments

In this section we present the results of some numerical experiments for the Galerkin procedure described above.

First, we study the sensitivity of the scheme to the Δt - H constraint. Consider (2.1) on the unit interval in \mathbf{R} with $a = 1$ and $b = 0$. Recall that, for purposes of L^2 -stability, the constraint is of the form

$$\Delta t < H^2/2. \quad (5.1)$$

We take as initial data the function u^0 given in Figure 1. At steady-state the solution $u(x, t) \equiv \int_0^1 u^0(x)$. We apply the Galerkin procedure (3.3), (3.6) to this problem, taking 4 subdomains. The L^2 norm of the solution for two different values of $\Delta t/H^2$ is given in Figure 2. Here it can be seen that when (5.1) is violated by as much as a factor of two, $\|U(\cdot, t)\|$ blows up as time increases.

Next, we study the experimental rate of convergence of the scheme. As remarked earlier, in certain instances a better rate of convergence than that predicted by Theorem 2 can be proven. However, for piecewise linear approximating spaces and $B(\psi)$ based on ϕ_k , $k = 2$ or 4, we expect the rate of convergence to be at best quadratic in h and order $k + 1$ in H , based on the truncation error of the scheme and the fact that the support of ϕ_k is $\mathcal{O}(H)$. Determining the sharpness of the estimates in general will be the subject of future work.

In these runs $a = 1$ and $b = 0$, and $\Omega = (0, 1) \times (0, 1)$. The solution u is approximated in the space of continuous piecewise linears in x tensored with continuous piecewise linears in y . We consider three scenarios:

1. Fully implicit Galerkin on uniform mesh; i.e., no domain decomposition,
2. Galerkin domain decomposition with four subdomains, with interfaces $\Gamma^i = \{i/4\} \times (0, 1)$, $i = 1, 2, 3$, and with global uniform mesh.
3. Galerkin domain decomposition with three subdomains Ω_i , $i = 1, 2, 3$, with interfaces

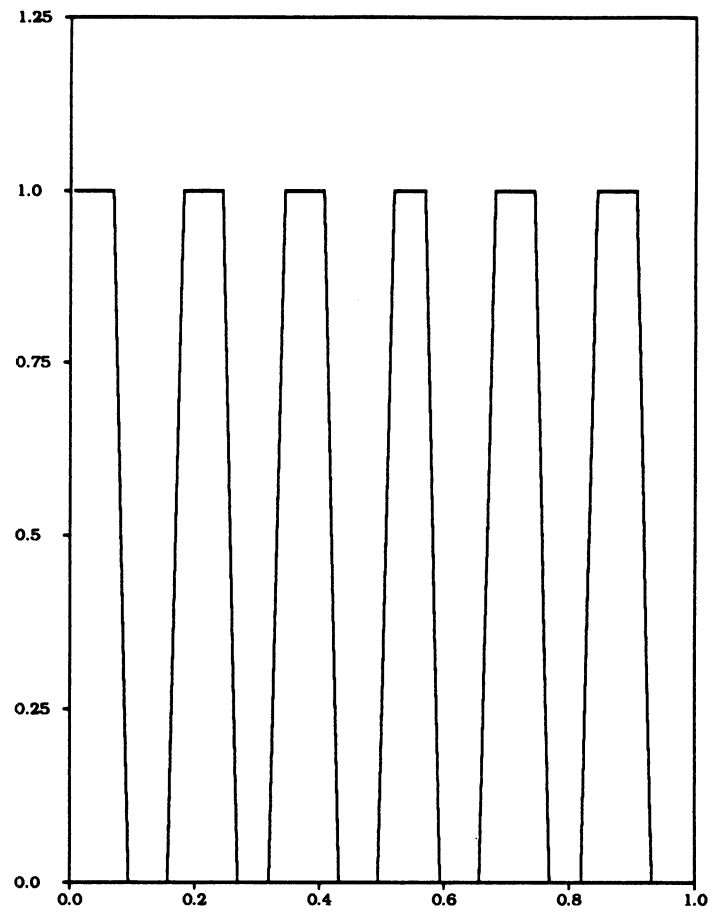


Figure 1: Initial data for stability test

